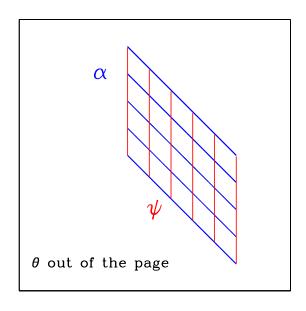
The Gyrokinetic Regime Geometry
Velocity Space
Linear How-To

#### **Coordinates**

 Given a magnetic field in Clebsch representation

$$\mathbf{B} = \nabla \alpha \times \nabla \psi$$

- Natural perpendicular coordinates are  $\psi$  and  $\alpha$ ; distance along the field line is  $\theta$
- ullet In general,  $(\psi, \alpha, \theta)$  is non-orthogonal
- Nor are they automatically single-valued.

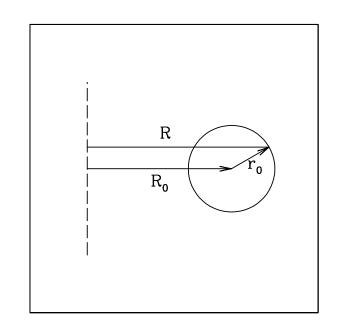


### **Simplest Toroidal Limit**

 Concentric circular flux surfaces:

$$\alpha = \phi - q(\theta - \frac{r}{R_0}\sin\theta)$$

$$\psi = \psi_0 + (r - r_0) \frac{B_0 r_0}{q}$$



• Note  $\theta$  is the distance along the field line.

(Concentric circles)

#### Flux Tube Limit

ullet Simulation domain is some number of gyroradii in  $\psi$  and  $\alpha$  directions, and  $\sim 2\pi qR$  along field line

Take advantage of ordering to write

$$h = \hat{h}(\theta) \exp(iS)$$

(Fluctuations vary slowly along field line, rapidly across)

Make contact with ballooning theory, defining

$$S = n_0 \left[ \alpha + q(\Psi) \theta_0 \right]$$

where  $n_0$  is the integer toroidal mode number.

### Theta Dependencies of $\nabla B$ Drift

- ullet Work out heta dependencies in various terms
- $\nabla B$  drift:

$$\frac{v_{\perp}^2}{2} \frac{1}{\Omega B^2} \mathbf{B} \times \nabla B \cdot \nabla S = \left(\frac{k_{\theta} \rho_i}{2}\right) \frac{v_{\perp}^2}{2} \left[\omega_{\nabla B} + \omega_{\nabla B}^{(0)} \boldsymbol{\theta}_{\mathbf{0}}\right],$$

where

$$\omega_{\nabla B} = \frac{2}{B^2} \frac{d\Psi}{d\rho} \hat{\mathbf{b}} \times \nabla B \cdot \nabla \alpha, \qquad \omega_{\nabla B}^{(0)} = \frac{2}{B^2} \frac{d\Psi}{d\rho} \hat{\mathbf{b}} \times \nabla B \cdot \nabla q.$$

• In the *high aspect ratio* concentric circles limit:

$$\omega_{\nabla B} = \frac{2a}{R_0} (\cos \theta + \hat{s}\theta \sin \theta) \qquad \omega_{\nabla B}^{(0)} = -\frac{2a}{R_0} (\hat{s}\sin \theta)$$

#### **Definitions**

• To get the previous expressions, define the normalized flux surface coordinate  $\rho=r/a$  and also  $k_{\theta}$ :

$$k_{\theta} \equiv \frac{n_0}{a} \frac{d\rho}{d\Psi_N} = \frac{n_0 q_0}{r_0}$$

- Normalized poloidal flux  $\Psi_N \equiv \Psi/(a^2B_a)$
- Here, a is half the diameter of the LCFS at the elevation of the magnetic axis.
- The field  $B_a$  is normalized to the toroidal field at  $R_a$ , where  $R_a$  is the avg of the min and max R on the LCFS.

## Theta Dependencies of Curvature Drift

Curvature drift:

$$\omega_{\kappa} = \omega_{\nabla B} + \frac{8\pi}{B^2} \frac{dp}{d\rho}.$$

where  $B = B(\theta)$  and p is the total pressure.

- In the high aspect ratio concentric circles limit  $\omega_{\nabla B} = \omega_{\kappa}$
- In general, the curvature drift is always *bad* on the outboard midplane of a tokamak, but the  $\nabla B$  drift can be reversed.

### $\theta$ Dependencies of Parallel Derivatives

• Parallel derivatives:

$$\hat{\mathbf{b}} \cdot \nabla = \frac{\mathbf{B}_0 \cdot \nabla}{B_0}$$

- In the *high aspect ratio* concentric circles limit  $\hat{\mathbf{b}} \cdot \nabla = a/(qR_0)$
- Freedom in the definition of  $\theta$  can be exploited to remove the  $\theta$  dependence from this operator (equal\_arc = T)

#### $\theta$ Dependencies of Perp Gradients

- Bessel functions have argument  $\gamma = |\nabla S|v_{\perp}/\Omega$ ; always enters as square  $(\gamma^2)$ .
- General expression:

$$|\nabla S|^2 = \frac{n_0^2}{a^2} |\nabla (\alpha + q\theta_0)|^2$$

$$= k_{\theta}^{2} \left( \frac{d\Psi}{d\rho} \right)^{2} \left| \nabla \alpha \cdot \nabla \alpha + 2\theta_{0} \nabla \alpha \cdot \nabla q + \theta_{0}^{2} \nabla q \cdot \nabla q \right|.$$

• In the high aspect ratio concentric circles limit

$$|\nabla S|^2 = k_\theta^2 \left| 1 + \hat{s}^2 \theta^2 - 2\theta_0^2 \hat{s}^2 + \theta_0^2 \hat{s}^2 \right|$$

### Flux-Surface Averages

 $\bullet$  Given Jacobian J, flux-surface average of a quantity  $\Gamma$  is

$$\langle \Gamma \rangle = \lim_{\Delta \rho \to 0} \frac{\int \Gamma J \, d\theta \, d\alpha \, d\rho}{\int J \, d\theta \, d\alpha \, d\rho}.$$

Flux-surface avg of a radially directed quantity (a flux) is

$$Q_{\mathsf{sim}} = \frac{\langle \mathbf{Q} \cdot \nabla \rho \rangle}{\langle \nabla \rho \rangle}$$

which will appear in the transport equation as

$$\frac{3}{2}\frac{d}{dt}\langle nT\rangle + \frac{1}{V'}\frac{d}{d\rho}AQ_{\text{sim}} + \dots = 0,$$

where the surface area  $A = 2\pi \langle |\nabla \rho| \rangle \int Jd\theta$ .

#### **Loose Ends**

Note the limit in this expression:

$$\langle \Gamma \rangle = \lim_{\Delta \rho \to 0} \frac{\int \Gamma J \, d\theta \, d\alpha \, d\rho}{\int J \, d\theta \, d\alpha \, d\rho}.$$

- In a flux-tube simulation, since the perpendicular box size is measured in gyroradii and  $\rho_* \ll 1$ , it is appropriate to take the integral over the entire simulation domain.
- In other words, a flux-tube simulation naturally calculates the flux-surface average of the various fluxes.

## Critical Elements for Axisymmetric Users

- Tokamak user has three choices:
  - 1. Assume shifted circles & specify r/a, a/R, q,  $\hat{s}$ , and  $\alpha$
  - 2. Specify the shape of the flux surface and  $B_p(\theta)$
  - 3. Read in numerical specification from MHD equilibrium solver
- In each case, the user can use the two free functions of the Grad-Shafranov equation to vary the magnetic shear and the pressure gradient freely: details controlled by bishop

# **Shifted Circles** (theta\_grid\_parameters)

- Magnetic shear:  $\frac{1}{s}$
- Aspect ratio:  $epsl = 2a/R_0$
- Safety factor:  $pk = 2a/(qR_0)$ ; i.e., q = eps1/pk
- Minor radius: eps = r/R
- Shafranov shift:  $shift = \alpha = -2Rq^2 d\beta/dr$  (not the coordinate  $\alpha$ !)

### Local Equilibrium

Surface of constant Ψ specified by:

$$R_N(\theta) = R_{0N}(\rho) + \rho \cos [\theta + \delta(\rho) \sin \theta],$$
$$Z_N(\theta) = \kappa(\rho)\rho \sin \theta.$$

(More general expressions easily implemented.)

Poloidal field determined from

$$B_p(\theta) = \frac{|\nabla \Psi|}{R} = \frac{d\Psi}{d\rho} \frac{|\nabla \rho|}{R}$$

• Conventions:  $R_N = R/a$ , etc.,

$$R_{0N}(\rho) = R_{0N}(\rho_c) + R'_{0N} d\rho, \quad \delta(\rho) = \delta(\rho_c) + \delta' d\rho, \quad \kappa(\rho) = \kappa(\rho_c) + \kappa' d\rho$$

### Local Equilibrium: Details I

• Input q, note that  $\oint \alpha \, d\theta = -2\pi q$  and then determine  $d\Psi/d\rho$  from

$$\frac{d\Psi}{d\rho} = \frac{I}{2\pi q} \oint \frac{d\theta}{R^2} (\nabla \theta \times \nabla \rho \cdot \nabla \phi)^{-1} = \frac{I}{2\pi q} \oint \frac{J \, d\theta}{R^2}$$

Define three integrals:

$$A(\theta) = \int J d\theta \left[ \frac{1}{R^2} + \left( \frac{I}{B_p R^2} \right)^2 \right], \qquad B(\theta) = I \int J d\theta \left[ \frac{1}{(B_p R)^2} \right],$$

$$C(\theta) = I \int J \, d\theta \left[ \frac{\sin u + R/R_c}{B_p R^4} \right]$$

### **Local Equilibrium: Definitions**

- $\bullet$   $u(\theta)$ : angle between the horizontal and the tangent to the magnetic surface in the poloidal plane
- $R_c(\theta)$ : local radius of curvature of surface in the poloidal plane.
- From papers by Mercier and Luc, used later by C. Bishop and R. Miller

## Local Equilibrium: Details II

• Upon defining  $\bar{A} = \oint d\theta \cdots$ , one can show

$$\hat{s} = \frac{\rho}{q} \frac{dq}{d\rho} = \frac{\rho}{2\pi q} \frac{d\Psi}{d\rho} \left( \bar{A}I' + \bar{B}p' + 2\bar{C} \right)$$

- Primes are derivatives w.r.t. Ψ
- User thus may specify any two of p', I', and  $\hat{s}$ . Several input options determined by bishop

### Local Eq Inputs (theta\_grid\_parameters)

• Magnetic shear:  $\frac{1}{shat} = \hat{s}$  (for bishop = 1)

• Minor radius:  $rhoc = \rho_c$  (Normalized by a!)

• Safety factor: qinp = q

• Elongation:  $akappa = \kappa$ 

• Deriv of elongation: akappri =  $d\kappa/d\rho$ 

### Local Eq Inputs (theta\_grid\_parameters)

• Triangularity:  $tri = \delta$ 

• Deriv of triangularity:  $tripri = d\delta/d\rho$ 

• Center of flux surf:  $Rmaj = R_{0N}$ 

• Local shift: shift =  $dR_{0N}/d\rho < 0$ 

• Center of LCFS:  $R_{geo} = R_{geoN}$  (for normalization only)

### Eq Inputs (theta\_grid\_eik\_knobs)

 These parameters may be set whether the equilibrium information comes from a local model or from numerical data (a file)

- Magnetic shear:  $s_hat_input = \hat{s}$  (for bishop > 1)
- Pressure gradient: beta\_prime\_input =  $d\beta/d\rho < 0$  (for bishop = 4)

### **Eq Selections** (theta\_grid\_eik\_knobs)

• Set at most one to true:

```
ppl_eq, transp_eq, efit_eq, gen_eq, vmom_eq, local_eq
```

- Set iflux = 1 to use numerical equilibrium, 0 otherwise
- Choose definition of  $\rho$  (flux surface label):

irho = 1 ... 
$$(\rho = \sqrt{\Phi/\Phi_a})$$
  
irho = 2 ...  $(\rho = d/D)$   
irho = 3 ...  $(\rho = \Psi/\Psi_a)$ 

Likely never change: itor = 1

### Eq Recommendations (theta\_grid\_eik\_knobs)

Watch out for the units, esp normalizations by a

Set writelots = T (it's essentially free information)

Set isym = 0 (allow up-down asymmetry)

Set equal\_arc = F (and avoid a rare bug)

Set del\_rho = 1.e-3 (nearly always adequate)

# Numerical Equilibria (theta\_grid\_eik\_knobs)

• If given the choice, use inverse solutions  $R(\Psi,\theta), Z(\Psi,\theta)$  for accuracy

Specify file with data using eqfile = '...'
 No blanks allowed, but other characters okay.

 Don't forget to set bishop. Use 1 to get the data from the file, > 1 to alter it.

#### References

- M. D. Kruskal and R. M. Kulsrud, Phys. Fluids, 1:265, 1958
- M. A. Beer and S. C. Cowley and G. W. Hammett, Phys. Plasmas, 2:2687, 1995
- R. L. Miller, et al., Phys. Plasmas, 5:973, 1998
- Notes for this talk: http://gk.umd.edu/g\_short.pdf